

Chapter VII: Key Results

Fermions

- Find $\{n_i\}_{mp}$ that maximizes $W_{FD}(\{n_i\}) = \prod_i \frac{g_i!}{n_i!(g_i - n_i)!}$ under the constraint $E = \sum_i n_i \epsilon_i = \text{constant}$

$$N = \sum_i n_i = \text{constant}$$

$$\frac{n_i}{g_i} = \frac{1}{e^\alpha e^{\beta \epsilon_i} + 1} = \frac{1}{e^{\beta(\epsilon_i - \mu)} + 1}$$

$$f_{FD}(\epsilon) = \frac{1}{e^{\beta(\epsilon - \mu)} + 1} \quad \text{Fermi-Dirac distribution}$$

Meaning: Number of fermions (particles) per single-particle state at energy ϵ

Bosons

- Find $\{n_i\}_{mp}$ that maximizes $W_{BE}(\{n_i\}) = \prod_i \frac{(g_i + n_i)!}{g_i! n_i!}$ under the constraint $E = \sum_i n_i \epsilon_i = \text{constant}$

$$N = \sum_i n_i = \text{constant}$$

$$\frac{n_i}{g_i} = \frac{1}{e^\alpha e^{\beta \epsilon_i} - 1} = \frac{1}{e^{\beta(\epsilon_i - \mu)} - 1}$$

$$f_{BE}(\epsilon) = \frac{1}{e^{\beta(\epsilon - \mu)} - 1} \quad \text{Bose-Einstein distribution}$$

Meaning: Number of bosons (particles) per single-particle state of energy ϵ

Using the constraint:

$$N = \sum_i g_i \frac{1}{e^{\beta(\epsilon_i - \mu)} + 1}$$

$$E = \sum_i g_i \frac{\epsilon_i}{e^{\beta(\epsilon_i - \mu)} + 1}$$

Look at the equation

Eq. for μ as a function of T

Using the constraint:

$$N = \sum_i g_i \frac{1}{e^{\beta(\epsilon_i - \mu)} - 1}$$

$$E = \sum_i g_i \frac{\epsilon_i}{e^{\beta(\epsilon_i - \mu)} - 1}$$

- Single-particle states are closed spaced in energy

\Rightarrow Continuous description $g(\epsilon)d\epsilon = \text{number of single particle states with energy in the interval } \epsilon \text{ to } \epsilon + d\epsilon$

$$N = \int \frac{g(\epsilon)}{e^{\beta(\epsilon - \mu)} + 1} d\epsilon$$

$$E = \int \frac{g(\epsilon) \epsilon}{e^{\beta(\epsilon - \mu)} + 1} d\epsilon$$

Eq. for μ as a function of T

$$N = \int \frac{g(\epsilon)}{e^{\beta(\epsilon - \mu)} - 1} d\epsilon$$

$$E = \int \frac{g(\epsilon) \epsilon}{e^{\beta(\epsilon - \mu)} - 1} d\epsilon$$

What's next?

$$N = \sum_i g_i \frac{1}{e^{\beta(\epsilon_i - \mu)} \pm 1}$$

$$E = \sum_i g_i \frac{\epsilon_i}{e^{\beta(\epsilon_i - \mu)} \pm 1}$$

Look into g_i

number of single-particle states
at some given energy

OR
more generally (practically)

$g(\epsilon)d\epsilon =$ number of single-particle states
in the energy interval ϵ to $\epsilon+d\epsilon$

(called "density of single-particle states"

OR "density of states")

$g(\epsilon)$ is NOT a Stat. Mech. problem. It is more a QM problem.

F. Discussions

- Here, we applied the Lagrange multipliers method as a quick and dirty way to obtain the Fermi-Dirac and Bose-Einstein distribution. Later, we will obtain these distributions again by invoking the Gibbs distribution and the Grand Canonical Ensemble, which are more convenient in treating quantum gases.
- Formally, the multipliers α and β should be determined by the constraints. Here, we claimed the established results of $\alpha = -\frac{\mu}{kT}$ and $\beta = \frac{1}{kT}$.
- The analysis also shows when quantum gases become classical ideal gas. We need $\frac{n_i}{g_i} \ll 1$, which amounts to separation $\gg \lambda_{\text{thermal}}$.
- We will treat ideal fermi gas and ideal bose gas in later chapters.
- The FD and BE distributions can also be obtained by minimizing F .

Summary

Students should be able to:

- carry out the process of minimizing or maximizing a quantity under some constraints, using the method of Lagrangian multipliers
- obtain the Fermi-Dirac and Bose-Einstein distributions
- understand what $f_{FD}(E)$ and $f_{BE}(E)$ mean
- understand that when $n_i \ll g_i$, the restrictions imposed by QM on single-particle state occupancy become un-important.
- use the Maxwell-Boltzmann distribution (after some practice) in classical statistical mechanics.

Refs:

The way (quick and dirty) we get the distributions is discussed in more elementary books.

- Tereza: Ch.3,4,5,6 ; Guénault: Ch.2,4,5
- All undergraduate books in Chinese.
 e.g. 苏汝璽, 「统计物理学」(第四章)
 黄昌德, 「热力学与统计物理学」(第七章)